

## Introduction

Supersymmetry, a symmetry between bosons and fermions, is one of the viable extensions of the Standard Model. When demanding that this symmetry is coordinate-dependent one automatically obtains a *supergravity* (SUGRA).

The mother of all supergravities is the 11-dimensional SUGRA. For some time now it has been conjectured [1] that the infinite-dimensional Lie algebra of the group  $E_{11}$  is its underlying symmetry.

Although D=11 SUGRA does not describe the world we live in, it is nonetheless interesting because it yields all maximal (ungauged) SUGRAS in lower dimensions upon compactification over a torus.

Gaugings can be introduced by either gauging directly in lower dimensions [2], or by flux compactifications of D=11 SUGRA. However, the latter is not able to reproduce all possible direct gaugings.

It has been shown that a particular decomposition of the low-level generators of  $E_{11}$  correspond to the physical fields of 11-dimensional supergravity [3], and other decompositions to the IIA and IIB SUGRAS in 10 dimensions [4].

This raises interesting questions ...

## Research questions

- Can the field content of all  $3 \leq D \leq 11$  maximal SUGRAS be obtained using various decompositions of  $E_{11}$ ?
- And if so, can some higher-level generators give rise to fields that do not propagate, but nonetheless deform (i.e. gauge) the theory?

## From Dynkin diagrams to Lie algebras

A compact way to represent a Lie algebra is the *Dynkin diagram*. It encodes the Cartan matrix  $A$ :

$$A_{ij} = \begin{cases} 2 & \text{if } i = j, \\ -1 & \text{if nodes } i \text{ and } j \text{ are connected,} \\ 0 & \text{otherwise.} \end{cases}$$

Assign to every  $i^{\text{th}}$  node three  $SL(2)$  generators:

- $h_i$  : Abelian (Cartan) subalgebra
- $e_i$  : Positive step operator
- $f_i$  : Negative step operator

They satisfy

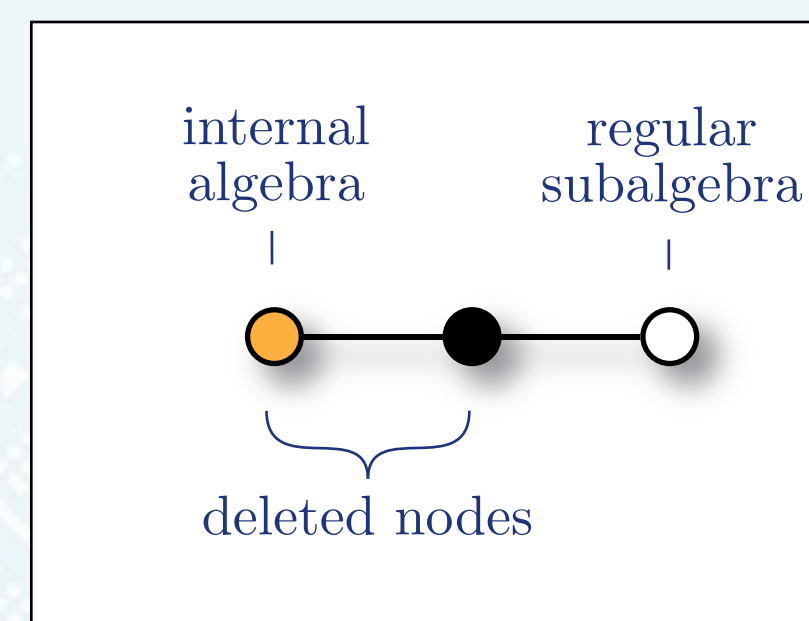
$$\begin{aligned} [h_i, h_j] &= 0, \\ [h_i, e_j] &= A_{ij} e_j, \\ [h_i, f_j] &= -A_{ij} f_j, \\ [e_i, f_j] &= \delta_{ij} h_i. \end{aligned}$$

The full Lie algebra is then generated by all multiple commutators of  $e$ 's and  $f$ 's, restricted by

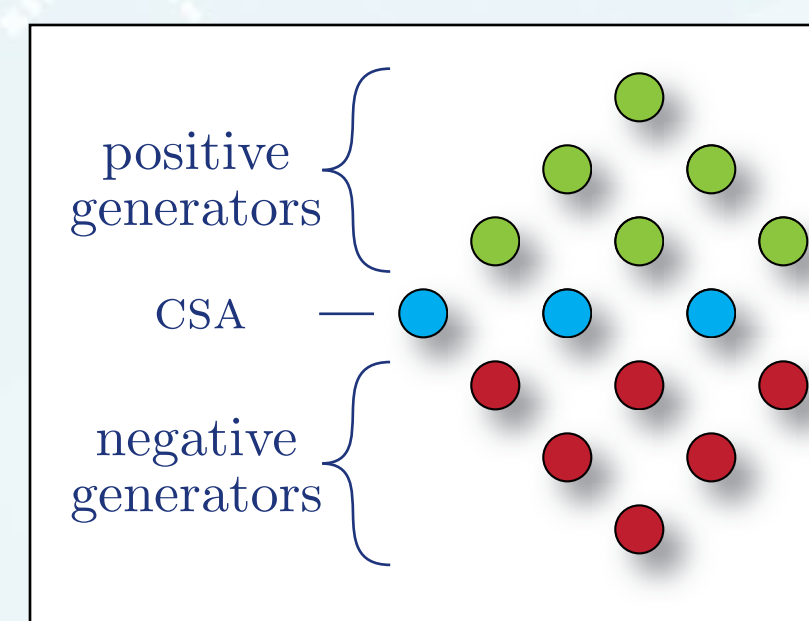
$$\begin{aligned} (\text{ad}_{e_i})^{1-A_{ij}} e_j &= 0, \\ (\text{ad}_{f_i})^{1-A_{ij}} f_j &= 0. \end{aligned}$$

## Easy example: $SL(4)$

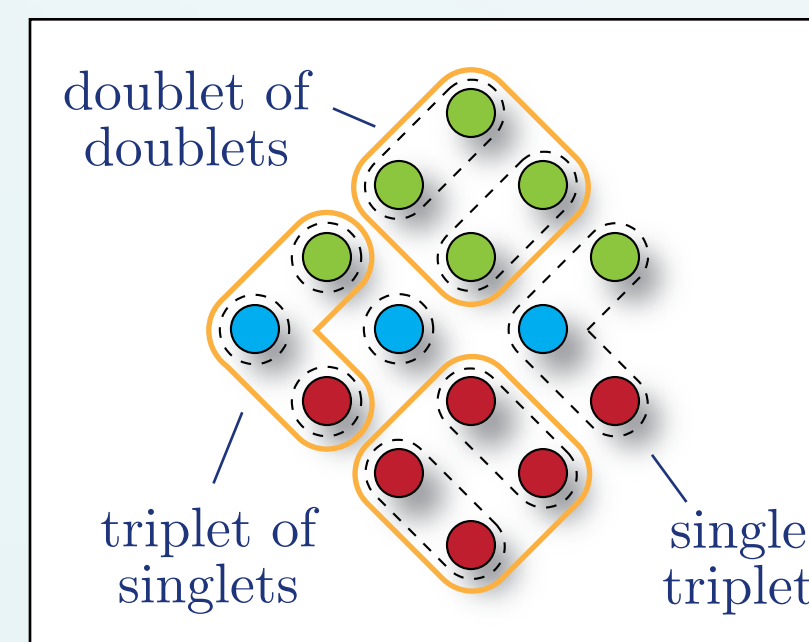
Two nodes in the diagram of  $SL(4)$  are 'deleted'. The node with no connections to the regular subalgebra forms the internal algebra.



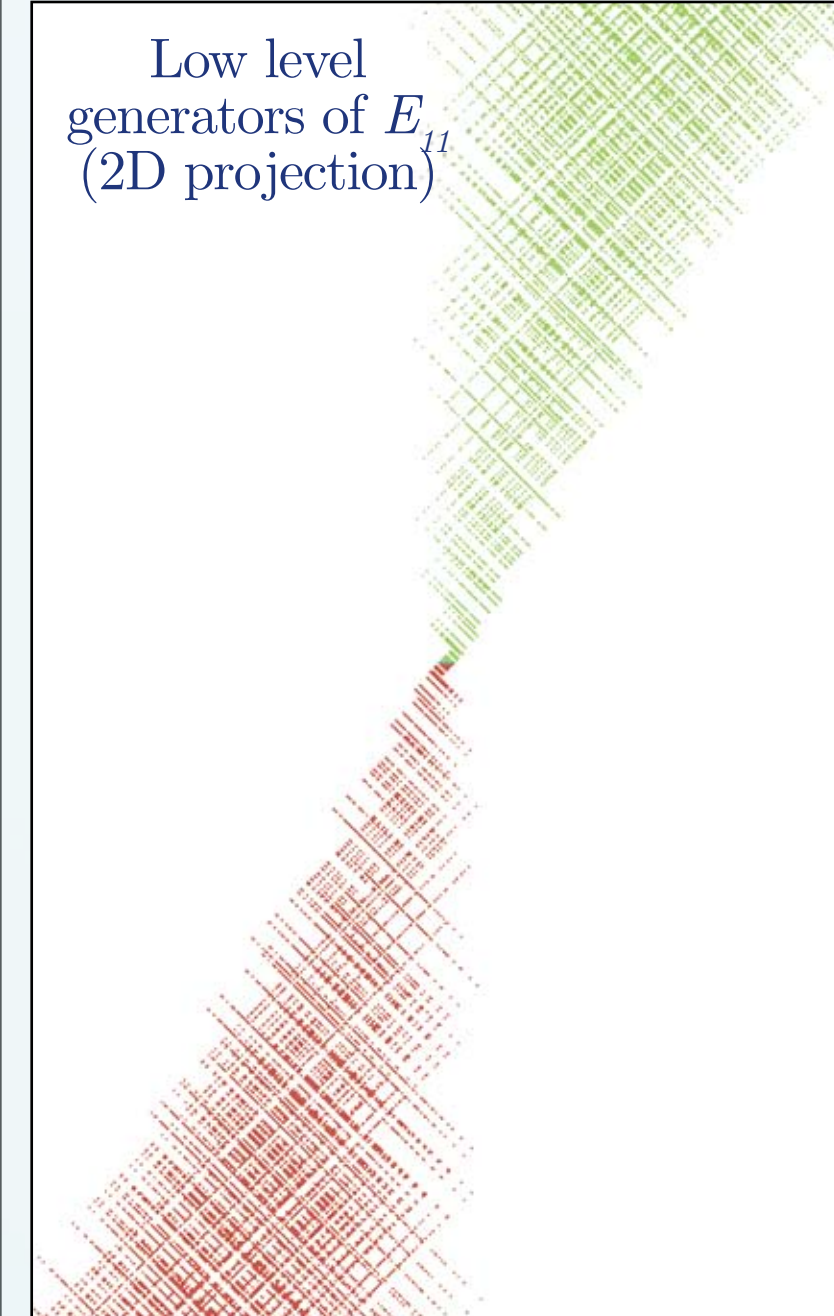
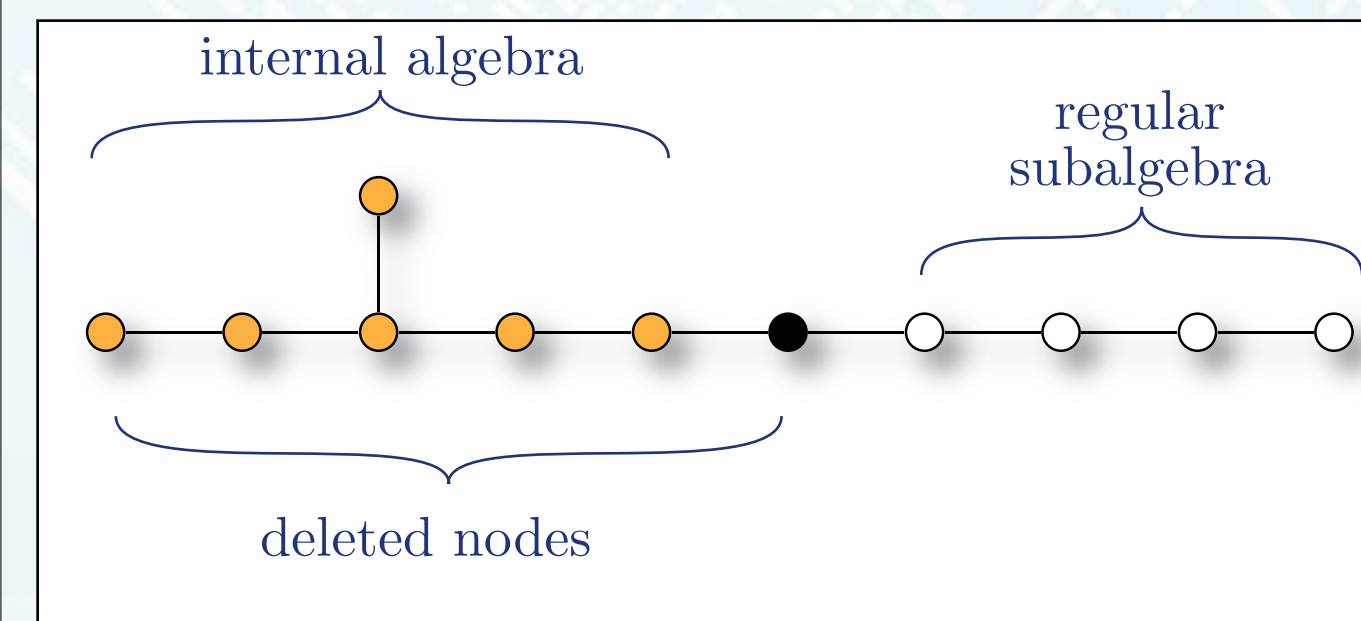
The Dynkin diagram gives rise to 15 generators: 6 positive, 6 negative, and 3 belonging to the Cartan subalgebra (CSA).



The 15 generators can be divided into representations of the regular (dotted lines) and internal (orange lines) subalgebra.



## The real deal: $E_{11}$



The Dynkin diagram of  $E_{11}$  can be decomposed in the same way. This gives rise to a similar ordering of the generators.

**Problem:**  
How do you order an infinite amount of generators?

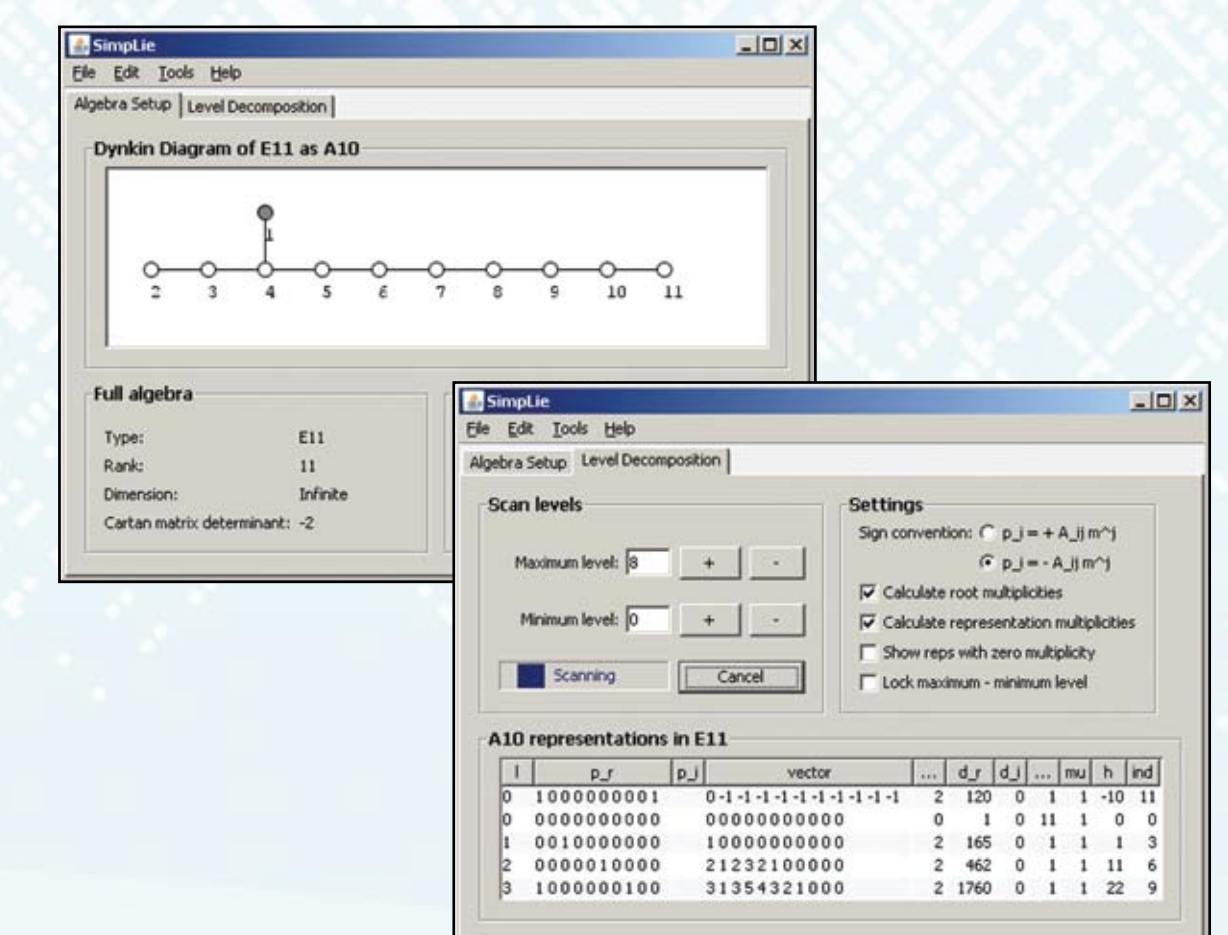
**Solution:**  
Let the computer figure it out!

## SimpLie

To speed up the decomposition of the  $E_{11}$  generators we have written a Java-based computer program, *SimpLie*.

From the user interface, one can:

- Create any simply-laced Dynkin diagram and 'delete' arbitrary nodes.
- Automatically scan for subalgebra representations at specified levels.



<http://strings.fmns.rug.nl/SimpLie>

## Results

Using SimpLie, we have found [5]:

- The field content of all  $3 \leq D \leq 11$  maximal SUGRAS as generators of  $E_{11}$ .
- Non-propagating 'de-form' and 'top-form' fields, which seem to have a close link to methods used to gauge these supergravities.

The de-forms and top-forms are not required by the supersymmetry algebra, and thus it cannot predict them. Our approach, and that of [6], shows that  $E_{11}$  can.